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# Science of tungstenbronze-type like $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$ (*R* = rare earth) microwave dielectric solid solutions

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# Abstract

The crystallography of tungstenbronze-type like  $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$  solid solutions for microwave dielectric materials applied for mobile phone and base station were studied. The crystal data and crystal structure of the fundamental lattice and supperlattice, and derivation of chemical formula and structural formula are summarized. Fourier map of the fundamental structure showed that oxygen ion positions of TiO<sub>6</sub> octahedra were splitted. This fact proved the tilting of octahedra which is cause of superlattice. The mechanisms of microwave dielectric properties are summarized. Finally, guidelines for the design of the microwave dielectric materials are presented. © 2001 Elsevier Science Ltd. All rights reserved.

Keywords: Ba<sub>6-3x</sub>R<sub>8+2x</sub>Ti<sub>18</sub>O<sub>54</sub> solid solutions; Internal strain; Microwave dielectric ceramics; Resonator; Tungstenbronze-type crystal structure

# 1. Introduction

The development of microwave telecommunication technology has been promoted by microwave dielectric ceramics. The tungstenbronze-type like  $Ba_{6-3x}R_{8+2x}$ . Ti<sub>18</sub>O<sub>54</sub> solid solutions as shown in Fig. 1 have been applied in mobile phones because of their high dielectric constant. This compound was reported by Kolar et al.<sup>1,2</sup> as  $BaO \cdot R_2O_3 \cdot 5TiO_2$  (R = rare earth) compound which was corrected to  $BaO \cdot R_2O_3 \cdot 4TiO_2$  later. About 10 years before Kolar's report, Bolton<sup>3</sup> had systematically studied the TiO<sub>2</sub>-rich region of the BaO–Nd<sub>2</sub>O<sub>3</sub>–TiO<sub>2</sub> ternary system. Following Kolar's work, Kawashima et al.<sup>4</sup> reported the compound in the R = Sm system and Nishigaki et al.<sup>5</sup> found another compound having good properties in the vicinity of 0.15(Ba<sub>0.95</sub>,Sr<sub>0.05</sub>)O/0.15Sm<sub>2</sub>O<sub>3</sub>/0.7TiO<sub>2</sub>.

Matveeva et al.<sup>6</sup> determined first the crystal structure of  $Ba_{3.75}Pr_{9.5}Ti_{18}O_{54}$  based on the fundamental unit cell by single crystal X-ray diffraction. This fundamental structure was confirmed by some researchers: Roth et al.<sup>7</sup> and Ohsato et al.<sup>8</sup> by single crystal X-ray diffraction, Kolar et al.<sup>9</sup> by Rietvelt method and so on. The superlattice with twice fundamental lattice spacing along the *c*-axis was pointed out by Matveeva et al.<sup>6</sup> and Kolar et al.<sup>9</sup> and was confirmed by an oscillation photograph of the single crystal<sup>10</sup> and by electron diffraction<sup>11</sup>. Ohsato et al.<sup>12</sup> and Rawn et al.<sup>13</sup>) reported the superstructure with  $TiO_6$  octahedral tilting. A different crystal structure with no tilting of  $TiO_6$  octahedra was presented as *Pbam* (No. 55) space group by Azough et al.<sup>11</sup>

This compound was found by Varfolomeev et al.<sup>14</sup> to form solid solutions as  $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$ . Ohsato et al.<sup>15</sup> also suggested the solid-solution formula and determined the solid solution regions for the Nd and Sm systems:  $0.0 \le x \le 0.7$  if R = Nd and  $0.3 \le x \le 0.7$  if R = Sm. The same solid solution region for the Sm system was also found by Negas et al.<sup>16</sup> With decrease in the ionic radius of R ion, the region of the solid solutions becomes narrower and for Gd and Eu, only  $BaO \cdot R_2$ - $O_3 \cdot 4TiO_2$  appeared.<sup>17</sup>

The dielectric properties as a function of composition x were reported for the Sm and Nd systems by Ohsato et al.<sup>18</sup> and Negas et al.,<sup>16</sup> and for the Pr system by Fukuda et al.<sup>19</sup> Ohsato et al.<sup>20</sup> found the special point with x = 2/3 at which the *Q*·*f* value becomes the highest due to the ordering in the rhombic and pentagonal sites. Valant et al.<sup>21</sup> studied the temperature coefficient of dielectric constants derived from the Clausius-Mosotti equation as a function of the ratio of mean radii ( $r_A/r_B$ ) of *A*- and *B*-site ions. Here,  $r_A/r_B$  is correlated to tilting of the TiO<sub>6</sub> octahedra.

In this study, we present crystal data and a crystal structure with tilting octahedra, and discuss the mechanisms of the microwave dielectric properties based on basic science, especially crystallography, and show the guideline for the design of the materials.

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## 2. Experimental

Crystal data were obtained with aid of single crystal X-ray diffraction using the Weissenberg camera and precession camera. The single crystals were grown by self-flux method. The superlattice was found using an oscillation photograph as reported in previous paper.<sup>10</sup>

Crystal structure analysis was performed by using a single crystal four-circle X-ray diffractometer. The  $\omega$ -2 $\theta$ scan technique was adopted to collect diffraction intensity data with graphite-monochromatized  $MoK_{\alpha}$  radiation. Three standard reflections showed no significant change during the measurement. The collection range for superlattice was  $2\theta < 91^{\circ}$ ,  $0 \le h \le 24$ ,  $0 \le k \le 44$ ,  $0 \le l \le 13$ . 8278 unique reflections were collected among which 6268 reflections with  $Fo > 3\sigma(|Fo|)$  were used for the structure analysis. The intensity data with l=2n for the fundamental lattice were selected from those of superlattice. The initial atomic parameters for the fundamental lattice were used from Matveeva et al.<sup>6</sup> and those of superlattice were derived from those of the fundamental lattice. The intensity data were corrected for Lorentz and polarization factors and absorption. Full-

Table 1 Crystal data of fundamental lattice and superlattice

matrix least-squares refinement on F of atomic parameters, with anisotropic thermal parameters and occupations of rare earth ions, was carried out with program RADY.<sup>22</sup>

The samples for the measurement of the microwave dielectric properties, the internal strain and lattice parameters were prepared by a solid state reaction as reported in previous papers.<sup>23</sup> The dielectric properties were measured by the Hakki and Colleman method.<sup>24</sup> The samples with optimal  $Q \cdot f$  values were selected for discussion.

The internal strain  $\eta$  was obtained from the following equation<sup>25</sup> as the grain size of the ceramics is sufficiently large:

 $\beta = 2\eta \tan \theta$ .

Here,  $\beta$  is the full-width at half-maximum (FWHM) of the X-ray powder diffraction peaks. The powder patterns were obtained by multi-detector system (MDS)<sup>26</sup> using synchrotron radiation in the "Photon Factory" of the National Laboratory for High Energy Physics in Tsukuba, Japan. The more precise measuring conditions were reported in previous paper.<sup>27</sup>

Accurate lattice parameters were obtained using the whole-powder-pattern-decomposition method (WPPD) program<sup>28</sup> for the powder diffraction patterns obtained by the step scanning technique using Si (99.99%) as an internal standard.

# 3. Crystal structure of $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$ solid solutions

# 3.1. Crystal data

Crystal data of both the fundamental lattice and superlattice are shown in Table 1. The naming of the axes is obeyed the rule b > a > c according to the conventional standard. The lattice spacing along the superlattice *c*-axis is twice that of the fundamental lattice. The morphology of the crystal is needle-like, elongated along the *c*-axis. Possible space groups of the fundamental lattice and the superlattice include *Pba2* (No. 32)

Fundamental lattice		Superlattice				
Crystal data		Crystal data				
Chemical formula	$Ba_{6-3x}Sm_{8+2x}Ti_{18}O_{54}(x=0.7)$	Chemical formula	$Ba_{6-3x}Sm_{8+2x}Ti_{18}O_{54}(x=0.7)$			
Crystal lattice type	Orthorhombic	Crystal type	Orthorhombic			
Space group	Pba 2(No. 32) or Pbam (No. 55)	Space group	<i>Pba</i> 2 <sub>1</sub> (No. 33) or <i>Pbnm</i> (No. 62)			
a (Å)	12.131(13)	a (Å)	12.131(13)			
$b(\dot{A})$	22.271(5)	$b(\dot{A})$	22.271(5)			
c (Å)	3.819(2)	c (Å)	7.639(5)			
Z	1	Z	2			
X-ray density (g/cm <sup>3</sup> )	5.91	X-ray density (g/cm <sup>3</sup> )	5.91			

or *Pbam* (No. 55), and *Pbn2*<sub>1</sub> (No. 33) or *Pbnm* (No. 62), respectively. As the result of crystal structural analysis, the space groups *Pbam* and *Pbnm* with center of symmetry were selected for the fundamental lattice and the superlattice, respectively. Recently, Ubic et al.<sup>29</sup> observed extra reciprocal lattice points, and suggested the space group may be  $Pb2_1m$  (No. 26), which lacks the *n*-glide plane perpendicular to *b*-axis.

# 3.2. Crystal structure

The final atomic parameters and Beq's of fundamental structure and superstructure are listed in Tables 2 and 3, respectively. The fundamental structure was converged with space group *Pbam*. The final reliability factor (R) is 4.49%. The fundamental structure is composed by three types of large cation sites: ten A1 rhombic sites in  $2 \times 2$ perovskite blocks, four A2 pentagonal sites and four trigonal sites. The pentagonal and trigonal sites are located among the perovskite blocks formed by four unit cells of perovskite. Fig. 2 shows a superimposed image of the electron density distribution of the fundamental structure and the framework of superstructure. We found splitting in the oxygen ion positions of  $TiO_6$ octahedra on the Fourier map. The splitting is due to tilting of TiO<sub>6</sub> octahedra as shown in Fig. 3. Therefore, the fundamental structure presents itself a mean structure of the real superstructure which octahedra are tilted along the *c*-axis. The change from the splitting to the tilting corresponds to the symmetry change from 2 to  $2_1$ 

Table 2						
Coordinations	on	the	fundamental	structure	of	Ba3.87Sm9.4Ti18O54
solid solution						

-						
Atoms	Site	g	X	У	Ζ	Beq. (Å)
Ti(1)	2a	1.0	0.5	0.5	0.0	0.23(9)
Ti(2)	4g	1.0	0.1983(2)	0.4353(1)	0.0	0.14(6)
Ti(3)	4g	1.0	0.3976(2)	0.1064(1)	0.0	0.27(6)
Ti(4)	4g	1.0	0.1175(2)	0.1634(1)	0.0	0.07(6)
Ti(5)	4g	1.0	0.3380(2)	0.2615(1)	0.0	0.09(6)
A 1(1)-Sm	4h	0.89	0.19986(9)	0.04815(5)	0.5	0.97(2)
A 1(2)-Sm	2d	0.92	0.0	0.5	0.5	0.91(3)
A 1(3)-Sm	4h	1.0	0.40687	0.37654(4)	0.5	0.22
A 2-Ba	4h	0.968	0.0902(1)	0.30644(6)	0.5	1.23(2)
O(1)	4h	1.0	0.100(2)	0.156(2)	0.5	5.4(3)
O(2)	4g	1.0	0.418(1)	0.1922(5)	0.0	0.3(3)
O(3)	4g	1.0	0.1837(9)	0.2371(6)	0.0	0.9(4)
O(4)	4h	1.0	0.333(2)	0.2788(6)	0.5	1.8(3)
O(5)	4g	1.0	0.377(1)	0.0193(6)	0.0	0.8(4)
O(6)	4h	1.0	0.211(3)	0.4424(8)	0.5	5.2(3)
O(7)	4g	1.0	0.242(1)	0.1119(6)	0.0	0.8(4)
O(8)	2b	1.0	0.5	0.5	0.5	10.0(6)
O(9)	4g	1.0	0.038(1)	0.0811(7)	0.0	6.2(2)
O(10)	4g	1.0	0.270(1)	0.3626(6)	0.0	0.9(4)
O(11)	4g	1.0	0.345(1)	0.4797(7)	0.0	4.2(1)
O(12)	4g	1.0	0.057(1)	0.4032(6)	0.0	4.2(1)
O(13)	4g	1.0	0.463(1)	0.3136(6)	0.0	1.6(6)
O(14)	4h	1.0	0.398(2)	0.0976(7)	0.5	3.7(3)

along the *c*-axis. Therefore, applying of space group  $P2_1/b2_1/a2/m$  (No. 55) to fundamental structure and  $P2_1/b2_1/n2_1/m$  (No. 62) to superstructure is considered to be more accurate and convincible. The obtained results, which point out the tilt of TiO<sub>6</sub> octahedra as a main factor for the superlattice formation, precludes the new crystal structure reported by Azough et al.<sup>11</sup> They applied space group *Pbam* (No. 55) to the superstructure. As Ti ions are located on the miller plane, the octahedra stand straight without tilting and with only a little rotation.

#### 3.3. Chemical formula

The distribution of cations in the crystal structure has an important impact on the dielectric properties. To understand more clearly the distribution of the cations, chemical formula and structural formulae are derived. Cations with large ionic radii are located within the three-dimensional framework of  $[\text{TiO}_6]^{8-}$  octahedra connected at all vertices. The atomic ratio of Ti to O is precisely 1:3 on the framework, and the perovskite-type structure also has this ratio. Then, in the BaO– $R_2O_3$ – TiO<sub>2</sub> (R = rare earth) ternary system, the solid solutions with the tungstenbronze-type like structure must be

Table 3

Coordinations on the superstructure of  $Ba_{3.87}Sm_{9.4}Ti_{18}O_{54}$  solid solution

Atom	Site	g	x	у	Ζ	Beq. (Å)
Ti(1)	4 <i>a</i>	0.5	0.0	0.0	0.0	0.38(5)
Ti(2)	8d	1.0	0.19861(11)	0.4352(6)	0.0005(2)	0.29(3)
Ti(3)	8d	1.0	0.39783(11)	0.10631(6)	0.0065(2)	0.41(3)
Ti(4)	8d	1.0	0.11712(10)	0.16348(5)	0.0026(2)	0.27(3)
Ti(5)	8d	1.0	0.33793(10)	0.26155(6)	0.0019(2)	0.31(3)
A 2(1)	4 <i>c</i>	0.494(1)	0.08528(7)	0.30018(3)	0.25	0.61(2)
A 2(2)	4 <i>c</i>	0.474	0.59587(7)	0.18661(4)	0.25	0.58(2)
A 1(1)	4c	0.396(1)	0.19193(6)	0.04409(3)	0.25	0.39(2)
A 1(2)	4 <i>c</i>	0.497(1)	0.70581(5)	0.44866(3)	0.25	0.39(2)
A 1(3)	4c	0.462(1)	0.99961(6)	0.49402(3)	0.25	0.33(2)
A 1(4)	4 <i>c</i>	0.501	0.40428(5)	0.37715(2)	0.25	0.34(2)
A 1(5)	4c	0.500	0.90971(5)	0.12416(3)	0.25	0.39(2)
O(1)	4 <i>c</i>	0.5	0.03996(90)	0.17214(37)	0.25	0.58(25)
O(2)	4 <i>c</i>	0.5	0.60537(72)	0.37016(35)	0.25	0.37(21)
O(3)	8d	1.0	0.41821(51)	0.19255(24)	0.0062(10)	0.57(16)
O(4)	8d	1.0	0.68511(49)	0.26288(26)	0.0219(8)	0.36(16)
O(5)	4 <i>c</i>	0.5	0.31213(79)	0.27916(40)	0.25	0.59(25)
O(6)	4 <i>c</i>	0.5	0.84572(78)	0.22324(42)	0.25	0.51(24)
O(7)	8d	1.0	0.37717(57)	0.01829(25)	0.0161(10)	0.70(18)
O(8)	4 <i>c</i>	0.5	0.23752(83)	0.44275(41)	0.25	0.61(26)
O(9)	4c	0.5	0.67319(84)	0.05498(40)	0.25	0.54(24)
O(10)	8d	1.0	0.24173(48)	0.11149(26)	0.0215(9)	0.40(15)
O(11)	4 <i>c</i>	0.5	0.48205(91)	0.4790(41)	0.25	0.66(27)
O(12)	8d	1.0	0.3748(59)	0.08057(28)	0.0487(11)	0.83(19)
O(13)	8d	1.0	0.76984(54)	0.13784(28)	0.0164(10)	0.71(18)
O(14)	8d	1.0	0.84388(52)	0.02003(31)	0.0376(10)	0.65(18)
O(15)	8d	1.0	0.05592(50)	0.40400(28)	0.0426(9)	0.51(16)
O(16)	8 <i>d</i>	1.0	0.46402(48)	0.31258(29)	0.0245(9)	0.52(17)
O(17)	4c	0.5	0.42941(104)	0.10066(42)	0.25	0.85(29)
O(18)	4 <i>c</i>	0.5	0.87883(85)	0.40424(45)	0.25	0.67(26)



Fig. 2. Electron density map (Fourier map) of the fundamental structure at z=0.5 revealing the top oxygen ion separated into two and superimposed a superstructure framework.



Fig. 3. TiO<sub>6</sub> octahedron tilting deduced from splitting of oxygen ion on the fundamental structure.

formed on the tie line between  $BaTiO_3-R_2Ti_3O_9$  where Ti/O ratio stays 1:3.  $BaO \cdot R_2O_3 \cdot 4TiO_2$  also locates on the tie line. Such solid solutions have different valence cations with large size: divalent  $Ba^{2+}$  and trivalent rare earth  $R^{3+}$ . To maintain electrostatic stability, three  $Ba^{2+}$  ions should be replaced with two  $R^{3+}$  ions and a vacancy. The substitutional formula of the solid solutions is considered as follows on this tie line.

$$3\mathrm{Ba}^{2+} \iff 2R^{3+} + V_{A1}$$

The derived chemical formula for the solid solutions is as follows:

$$Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$$
.

Here, the formula  $Ba_6R_8Ti_{18}O_{54}$  is obtained as the end member of the solid solutions, in which the large

sites: A1 rhombic sites and A2 pentagonal sites (more detailed description of sites is presented in the previous section) are fully occupied. There is one formula unit in the fundamental unit cell which contains 18 Ti<sup>4+</sup> ions and 54 O<sup>2-</sup> ions. The BaO· $R_2$ O<sub>3</sub>·4TiO<sub>2</sub> composition corresponds exactly to x=0.5, and Ba<sub>3.75</sub>Pr<sub>9.5</sub>Ti<sub>18</sub>O<sub>54</sub> to x=0.75.

There are three different cations with different diameter in the crystal structure, which occupy different sites. The middle-sized R ions mainly occupy the A1-sites and the largest Ba ions mainly occupy pentagonal A2-sites. If the composition is Ba rich, small amount of Ba ions occupy also A1-sites. The smallest Ti ions alone occupy octahedra B-sites. The trigonal C-sites are empty.

The solid solution regions are  $0.0 \le x \le 0.7$  for R = La, Pr and Nd, and  $0.3 \le x \le 0.7$  for R = Sm. In the La, Pr and Nd system, the substitution mode of cations undergoes some changes.

# 3.4. Structural formula

The basic structural formula is as follows:

# $[S_{10}]_{A1}[S_4]_{A2}B_{18}X_{54}$

Ten S sites in A1 and four S sites in A2 are included in a fundamental unit cell. Here, B is a cation in an octahedron and X is an anion, occupied by Ti and O ions, respectively. The final composition  $Ba_6R_8Ti_{18}O_{54}$  of the solid solutions with x = 0 is represented as the structural formula  $[R_8Ba_2]_{A1}[Ba_4]_{A2}Ti_{18}O_{54}$  in which the two largest Ba ions together with eight middle-sized R ions occupy A1-sites. The A2-sites are occupied only by four Ba ions. The structural formula of the solid solution is derived as follows from the terminal composition according to the substitution mode  $3Ba^{2+} \iff 2R^{3+} + V_{A1}$ .

 $[R_{8+2x}Ba_{2-3x}V_x]_{A1}[Ba_4]_{A2}Ti_{18}O_{54}$ in the range of  $0 \le x \le 2/3$ 

When all the Ba ions in A1-sites are substituted by R ions, x=2/3. Moreover, for  $x \ge 2/3$ , Ba ions in A2-sites become substituted by R ions, which occupy the A1-sites, and vacancies can be created in A2-sites. The structural formula is as follows:

 $[R_{9+1/3+2(x-2/3)}V_{2/3-2(x-2/3)}]_{A1}[Ba_{4-3(x-2/3)}V_{3(x-2/3)}]_{A2}$ Ti<sub>18</sub>O<sub>54</sub> in the range of 2/3  $\leq x \leq 1$ .

The structural formulae derived above, divide the large cation sites into two kinds: rhombic and pentagonal sites. In the more precise fundamental structure, the A1 site separates into three sites as follows:

 $[S_2]_{A11}[S_2]_{A12}[S_2]_{A13}[S_2]_{A2}B_{18}X_{54},$ 

and in the real superstructure, A1 sites separate into five sites, and A2 sites separate into two sites as follows:

 $[S_2]_{A11}[S_2]_{A12}[S_2]_{S13}[S_2]_{A14}[S_2]_{A15}[S_2]_{A21}[S_2]_{A22}B_{18}X_{54}.$ 

The site occupancies obtained at x=0.7 are represented as follows:

$$\begin{split} & [Sm_{1.58}]_{\mathcal{A}11} [Sm_{1.99}]_{\mathcal{A}12} [Sm_{1.85}]_{\mathcal{A}13} [Sm_{2.00}]_{\mathcal{A}14} [Sm_{2.00}]_{\mathcal{A}15} \\ & [Ba_{1.98}]_{\mathcal{A}21} [Ba_{1.90}]_{\mathcal{A}22} Ti_{18} O_{54} \end{split}$$

In the case of x = 0.5 the precise structural formula is as follows:

$$\begin{split} & [Sm_{1.68}]_{\mathcal{A}11}[Sm_{1.96}]_{\mathcal{A}12}[Sm_{1.88}]_{\mathcal{A}13}[Sm_{1.68}Ba_{0.31}]_{\mathcal{A}14} \\ & [Sm_{1.80}Ba_{0.18}]_{\mathcal{A}15}[Ba_{2.00}]_{\mathcal{A}21}[Ba_{2.00}]_{\mathcal{A}22}Ti_{18}O_{54} \end{split}$$

The A14 and A15 sites are occupied by Sm and Ba ions disorderly. As the determination of the site occupancy is difficult, some researchers reported different results as follows:

 $[Sm_{1.72}]_{A11}[Sm_{2.00}]_{A12}[Sm_{0.81}Ba_{1.19}]_{A13}[Sm_{2.00}]_{A14}$  $[Sm_{2.00}]_{A15}[Ba_{2.00}]_{A21}[Ba_{2.00}]_{A22}Ti_{18}O_{54}$  $x = 0.27 \text{ by Rawn et al.}^{13}$ 

and

 $[Sm_{1.44}Ba_{0.24}]_{A11}[Sm_{1.99}]_{A12}[Sm_{1.65}Ba_{0.26}]_{A13}[Sm_{1.94}]_{A14}$  $[Sm_{2.00}]_{A15}[Ba_{2.00}]_{A21}[Ba_{2.00}]_{A22}Ti_{18}O_{54}$  $x = 0.5 \text{ by Okudera et al.}^{30}$ 

#### 4. Properties and structure

# 4.1. High quality factor due to ordering of Ba and R

The dielectric properties as a function of the composition x for the  $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$  solid solutions<sup>18,20</sup> are shown in Fig. 4. Similar results for Sm and Nd system are shown by Negas et al.<sup>16</sup> The characteristic phenomenon is that  $Q \cdot f$  values varied non linearly as a function of composition, though  $\varepsilon_r$  and  $\tau_f$  vary proportional to the composition. The data of the x = 2/3 composition in which R and Ba ions occupy separately the rhombic sites (A1) and the pentagonal sites (A2), respectively, show the highest Q.f values: 10549 GHz in the Sm system, 10010 GHz in the Nd system and 2024 GHz in the La system as shown in Fig. 4. The highest quality factor may be based on the ordering of R and Ba ions in the A1 and A2 sites, respectively. The distribution of the ions might reduce the internal strain and result in the non-linear variation of quality factor.

Internal strain  $\eta$  values for x = 0.3, 0.5, 2/3 and 0.7 are shown in Fig. 5. It should be noticed that the internal strain for x = 2/3 is the lowest in the series of Ba<sub>6-3x</sub> Sm<sub>8+2x</sub>Ti<sub>18</sub>O<sub>54</sub> solid solutions. This low internal strain comes from the distribution of cations in the rhombic sites and the pentagonal sites on the tungstenbronze-



Fig. 4. Dielectric properties of the tungstenbronze-type like  $Ba_{6-3x}R_{8+2x}Ti_{18}O_{54}$  solid solutions as a function of composition.



Fig. 5. Internal strain obtained from the equation  $\beta = 2\eta \tan\theta$  using FWHM of powder XRD (a), and internal strain vs. composition (b).

type like structure. In the x=2/3 composition, ions with the same size occupy each A1 and A2 site, as shown in the structural formula  $[R_{9.33}V_{0.67}]_{A1}[Ba_4]_{A2}Ti_{18}O_{54}$ , that means, R ions and Ba ions are ordering in both the rhombic sites and pentagonal sites, respectively. This ordering leads to the lowest strain. As the x-values decrease according to the structural formula  $[R_{8+2x}Ba_{2-3x}V_x]_{A1}[Ba_4]_{A2}Ti_{18}O_{54}$  in the range of  $0 \le x \le 2/3$ , Ba ions with larger ionic radii will occupy also a part of the rhombic sites with their smaller size. The occupation of Ba ions in A1 site leads to internal strain around themselves which lowers the  $Q \cdot f$  values. Moreover, the vacancies generated in A1 sites by the substitution of 3Ba by 2R might be the second reason for lowering the internal strain and lead to the high  $Q \cdot f$  values. On the other hand, as the x-values increase according to the structural formula  $[R_{9.33+2(x-2/3)}V_{0.66-(x-2/3)}]_{A1}[Ba_{4-3(x-2/3)}]_{A1}$  $V_{3(x-2/3)}]_{A2}$ Ti<sub>18</sub> O<sub>54</sub> in the range of  $2/3 \le x \le 0.7$ , Ba ions in A2 sites are substituted by R ions. The decrease of Ba ions produces vacancies in A2 sites and may lead to unstable crystal structures, as shown by the limit of solid solubility located near the x=0.7 composition. Moreover, the decrease in the number of vacancies in the rhombic sites accompanied by the decrease of Ba ions in the pentagonal sites might lead to an additional internal strain. These strains are the reason for the lower quality factor at the x = 0.7.

On the other hand, the  $Q \cdot f$  values of each *R*-compound with x = 2/3 in the Ba<sub>6-3x</sub> $R_{8+2x}$ Ti<sub>18</sub>O<sub>54</sub> solid solutions increase according to a decrease in the rareearth ion size (lanthanoid contraction). The Sm-compound has a better  $Q \cdot f$  than the La-compound. This crystal structure is maintained by the size difference of large cations such as Ba and *R*. It was revealed that the crystal structure with the largest size difference between Ba and *R* shows excellent quality factor, as it has low internal strain.

# 4.2. Dielectric constant and temperature coefficient of resonant frequency

The dielectric constant ( $\varepsilon_r$ ) is affected by the following three effects: (I) volume of TiO<sub>6</sub> octahedra, (II) tilting of

Table 4 Volume and tilt angles of  $TiO_6$  octahedra strings along the *c*-axis about two compositions of x = 0.5 and 0.7

	х	Ti(1)	Ti(2)	Ti(3)	Ti(4)	Ti(5)	Mean
Volume (Å <sup>3</sup> )	0.5	9.864	10.04	9.791	9.737	10.30	9.946
	0.7	9.858	10.04	9.777	9.761	10.19	9.925
Tilt angle $(\theta)$	0.5	14.45	10.80	8.996	10.03	5.672	9.990
	0.7	15.10	11.56	9.261	11.07	6.161	10.63

octahedra strings and (III) polarizabilities of *R* and Ba ions. The dielectric constants of the solid solutions are proportional to lattice parameters or cell volumes as shown in Fig. 7a. As *x* increased,  $\varepsilon_r$  decreased linearly and lattice parameters or cell volumes also decreased linearly.<sup>20</sup>

Usually, in the perovskite structure, the polarity of the Ti ion in the octahedra is produced as a result of the large octahedral volume. Thus, as the mean value of the volume decreased from 9.946 Å<sup>3</sup> at x=0.5 to 9.925 Å<sup>3</sup> at x=0.7 as shown in Table 4, this volume change is considered to have decreased  $\varepsilon_r$ . However, the volume change is very small, thus other effects should be examined, such as tilting of the TiO<sub>6</sub> octahedra strings as suggested by Valant et al.<sup>21</sup> The tilting angle, which is between the *c*-axis and the central axis of the octahedra as shown in Fig. 6, is inverse proportional to lattice parameters: the mean tilting angle is 9.99° at x=0.5 and 10.65° at x=0.7, based on the refined crystal structure of the Sm solid solution series as shown in Table 4.

 $\varepsilon_{\rm r}$  and  $\tau_{\rm f}$  for the solid solution regions with  $0.3 \le x \le 0.7$  for the Sm system,  $0 \le x \le 0.7$  for the Nd, Pr and La systems are replotted as a function of unit cell volume in Fig. 7. These  $\varepsilon_{\rm r}$  and  $\tau_{\rm f}$  values are almost proportional to the unit cell volume. The size of cell volume affects to the  $\varepsilon_{\rm r}$  such as large size brings large  $\varepsilon_{\rm r}$ . From



Fig. 6. Dielectric constant  $\varepsilon_r$  concerned with tilting angle.

this figure, it was also deduced that the polarizabilities of R ions affect  $\varepsilon_r$  and  $\tau_f$ . On the table of the polarizabilities derived by Shannon,<sup>31</sup> the La ion, which gives the largest  $\varepsilon_r$  in the series, also has the largest polarisability among these R ions: 6.03 for La, 5.31 for Pr, 5.01 for Nd and 4.74 Å<sup>3</sup> for Sm. The  $\varepsilon_r$  values decrease with the polarizabilities. On the other hand, the  $\varepsilon_r$  values also vary linearly as a function of cell volume in each R-system as shown in Fig. 7a. The variations of  $\varepsilon_r$  are also affected by the polarizabilities of R and Ba ions. The substitution is performed according to the following equation:  $3Ba \rightarrow 2R + V$ . The total polarizabilities due to the substitution equation are reduced from  $3 \times 6.40$  to  $2 \times 6.03$  Å<sup>3</sup> for the La system. Here, the value of 6.40 Å<sup>3</sup> for the Ba<sup>+2</sup> ion is larger than that for the La<sup>+3</sup> ion.

The  $\tau_f$  are also plotted as a function of cell volume in Fig. 7b. Though similar tendency with  $\varepsilon_r$  is observed, the mechanism of  $\tau_f$  has not been clarified yet. The  $\tau_f$  values of the Sm system are usually negative but close to zero. As  $\tau_f$  obeys additional rule, we could get easily a material with the  $\tau_f=0$  ppm/°C. In the Sm system adding Nd or La system, outstanding materials with  $\tau_f=0$  ppm/°C have been realized by Ohsato.<sup>32</sup>

# 5. Guidelines for the design of the materials

1. Knowledge of the crystal structure is basically very important, such as super structure, ordering and distribution of cations, for the development of high  $Q \cdot f$ 



Fig. 7. The dielectric constant (a) and the temperature coefficients of resonant frequency (b) as a function of the unit cell volume.

ceramics. 2. Selection of the structure without internal strain for high Q (low loss) dielectric materials. The internal strain is reduced by ordering of cations. In this case, ordering composition of Ba and R ions at x=2/3brings low internal strain and high Q. 3. Crystal structure with centrosymmetry is better for high  $Q \cdot f$ . 4. For high  $\varepsilon_r$ or high polarity, selection of large octahedral volume, and small tilting angle. These factors are found for these solid solution series. 5. For high  $\varepsilon_r$ , selection of atoms with large polarizability values proposed by Shannon.<sup>31</sup> The  $\varepsilon_r$  increases with polarizability: Eu (4.53) $\Rightarrow$ Sm  $(4.74) \Rightarrow \text{Nd} (5.01) \Rightarrow \text{Pr} (5.31) \Rightarrow (6.03) \Rightarrow \text{Ba}(6.40) \text{ A}^3. 6.$ Value of  $\tau_f = \text{zero ppm}/^{\circ}C$  could be achieved by additive rule using the materials with positive and negative  $\tau_{\rm f}$  or making solid solutions between composition with positive and negative  $\tau_{\rm f}$ . The example is shown in the Sm– Nd and Sm-La systems.<sup>32</sup>

## 6. Summary

The space groups, which were presented by many researchers, were confirmed as Pbam for fundamental lattice and *Pbnm* for superlattice. The crystal structure is refined as tungstenbronze-type like structure with  $2 \times 2$ perovskite blocks by X-ray crystal structure analysis. The superstructure is caused by the tilting of octahedra strings along the *c*-axis which was deduced from splitting of oxygen ion position on the Fourier map of the fundamental lattice. The chemical and structural formulae were derived. Mechanisms of dielectric properties are stated, and ordering of Ba and R ions at x=2/3composition is considered to result in high quality factors. The dielectric constants are found to be affected by three kinds of factors: volume of octahedron, tilting of octahedron and polarizabilities of ions. Finally, guidelines for the design of dielectric materials are suggested.

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